

Laser Excitation of the Thorium-229 Nucleus – Towards a Nuclear Clock

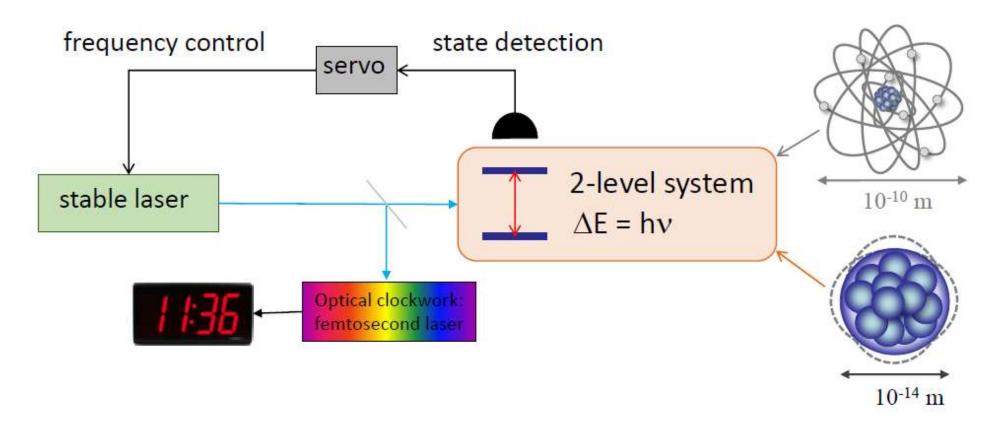
Ekkehard Peik

Time and Frequency Department PTB, Braunschweig, Germany



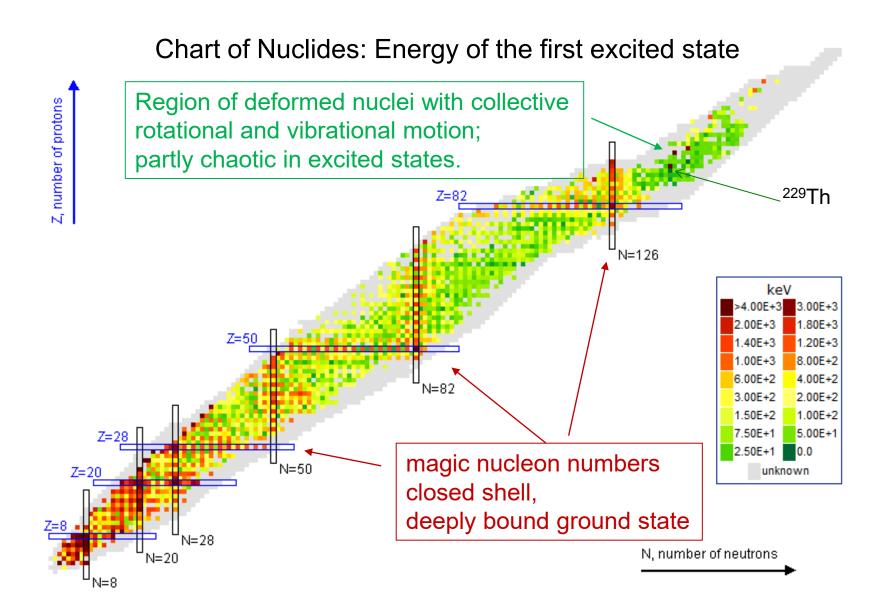
From the atomic to the nuclear clock





The nuclear clock promises:

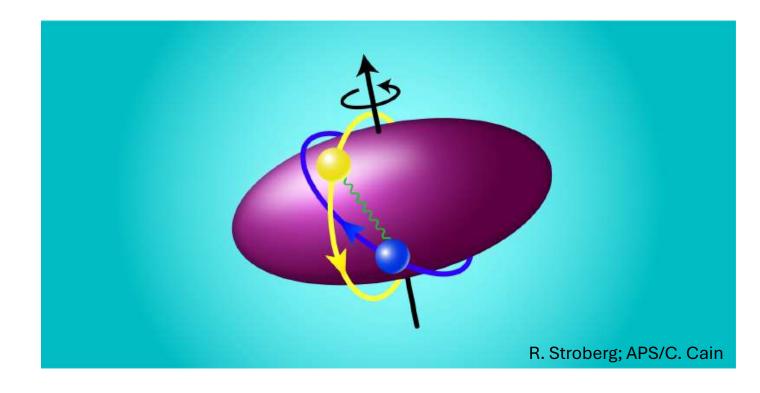
- High accuracy (with laser cooled trapped ions)
- High stability (in the solid state as a laser Mössbauer system)
- High sensitivity to new physics (also to strong interaction)

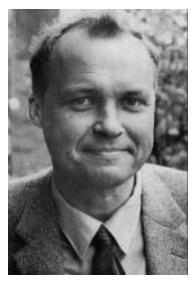


From: National Nuclear Data Center, Brookhaven National Laboratory

Nuclear structure: The deformed shell model or Nilsson model

Collective motion of an elliptical core combined with single-particle motion of an unpaired "valence" nucleon (much faster than motion of the core)





Sven Gösta Nilsson

Not a complete description of the Th-229 transition:

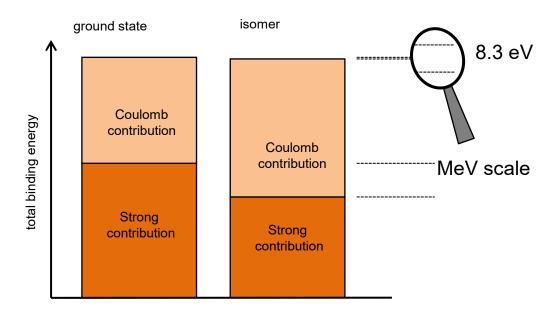
The proton distribution in the core would not change.

But we have measured a change in the rms charge radius:

$$\langle r_{229\text{m}}^2 \rangle - \langle r_{229}^2 \rangle = 0.0105(13) \text{ fm}^2$$
 J. Thielking et al., Nature **556**, 321 (2018)

How does the very low energy appear?

"Magic" cancellation of contributions from Coulomb and strong forces V.V. Flambaum, Phys. Rev. Lett. 97, 092502 (2006)



→ High sensitivity of a Th-229 nuclear clock in fundamental tests: Search for variations of fundamental constants or other violations of the Einstein equivalence principle

Nuclear Clock:

Oscillator that is frequency-stabilized to a nuclear (γ -ray) transition

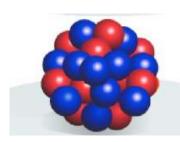
Higher accuracy:

The nuclear clock allows for a choice of a suitable electronic state for the interrogation of the nuclear resonance, for example in laser-cooled ions, Th³⁺.

Example: Electric fields, Stark effect:



Electrons shield the nucleus (from homogeneous, static fields)



Nuclear polarizability is orders of magnitude smaller; electronic effects are common-mode for both states of the transition.

Analyses of best suited electronic configurations:

E. Peik, Chr. Tamm, Europhys. Lett. **61**, 181 (2003)

C. J. Campbell et al., PRL 108, 120802 (2012)

→ low values of electronic angular momenta J=0 or J=1/2

→ stretched states F=J+I, aligned electronic and nuclear momenta

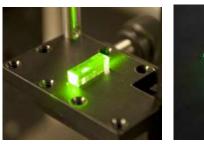
Nuclear Clock:

Oscillator that is frequency-stabilized to a nuclear (γ -ray) transition



Higher stability:

In a Mössbauer solid state nuclear clock, many absorbers may be interrogated (>10¹⁵ instead of ≈10⁰ (ion trap) or ≈10⁴ (optical lattice)). Systematics: Crystal field shifts. Proposed at PTB, UCLA, TU Wien



UCLA, LiCAF

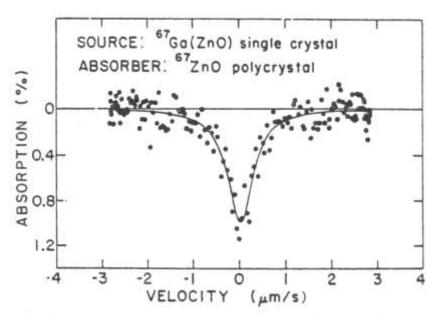
Britania anna

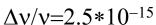
TU Vienna, CaF₂

Nuclear gamma spectroscopy provides very high spectral resolution

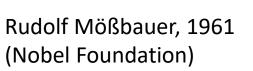
Mössbauer effect: Recoil-free resonance absorption of gamma radiation by nuclei in a crystal lattice

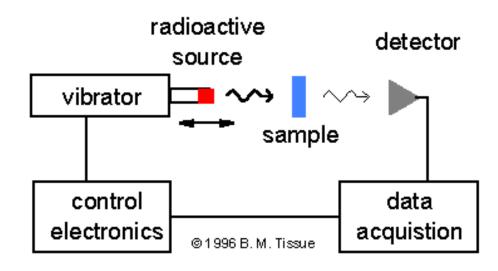
Mößbauer spectrum of the 93.3 keV resonance of Zn-67





G. Perlow et al., 1974





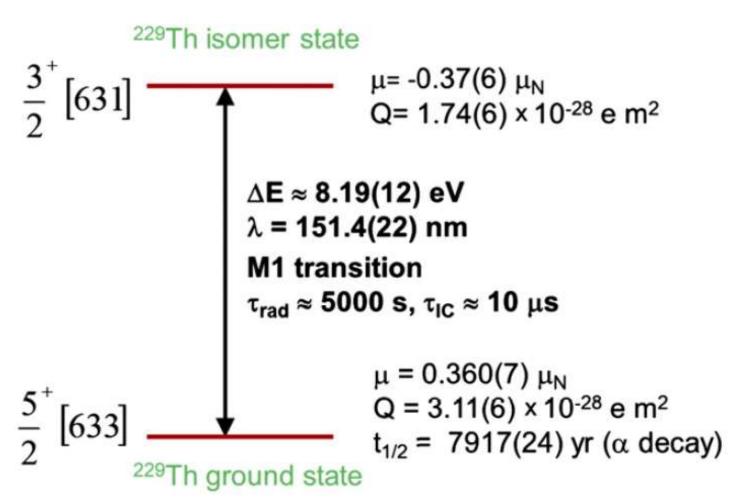
Nuclear Clock:

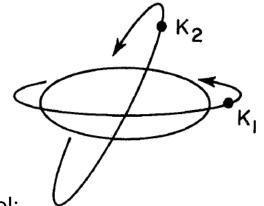
Replace the incoherent radioactive source with a tunable laser



The Th-229 low-energy isomer: State of knowledge in 2020

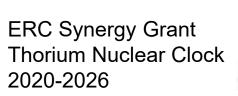






Nilsson Model:
Deformed core plus unpaired neutron

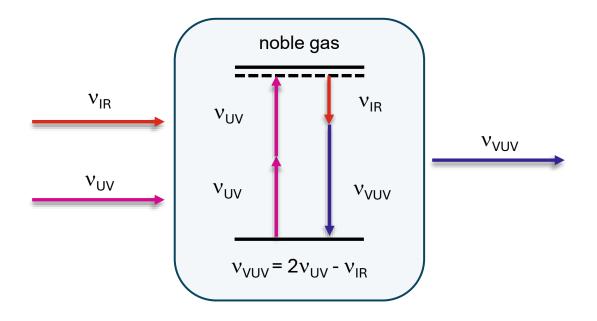
E. Peik, T. Schumm, M. Safronova, A. Pálffy, J. Weitenberg, P.G. Thirolf, Quant. Sci. Tech. 6, 034002 (2021)





8-eV tunable VUV generation: Four wave mixing

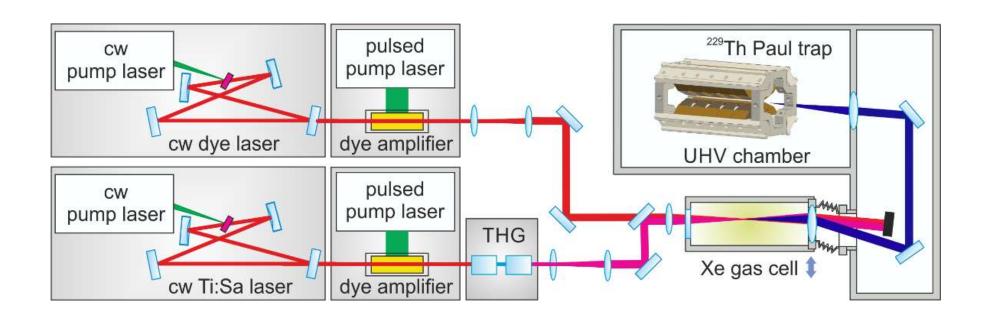
- Near-resonantly driving of 2-photon transition in noble gas.
- Supplying one photon with v_{IR} yields fourth photon with $v_{VUV} = 2v_{UV} v_{IR}$.
- Two-photon transition in Xe at 2 x 250 nm is suitable for VUV tunability from 167 nm to 148 nm, i.e. 7.42 eV to 8.38 eV.

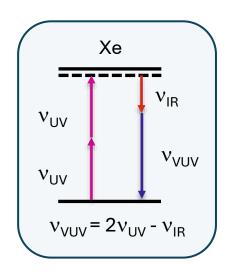


See e.g.: S. J. Hanna et al., Int. J. Mass. Spectrom. **279**,134 (2009)

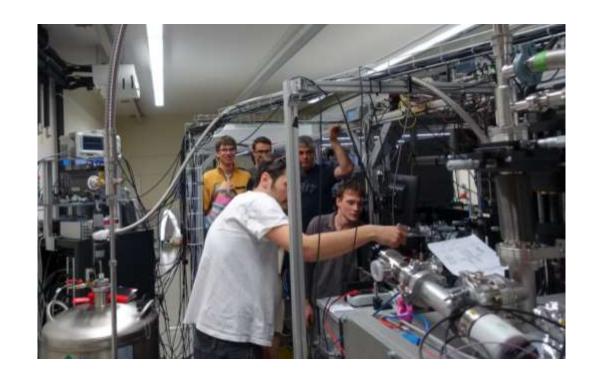
8-eV tunable VUV generation: Four wave mixing

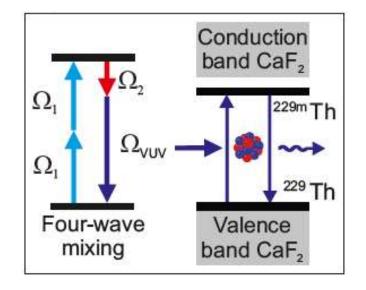
- Generating 148 nm VUV requires laser beams at 250 nm and 790 nm.
- Third order process needs high intensity to achieve suitable efficiency.
- Pulsed lasers (~10 ns, 30 Hz repetition rate) best compromise between VUV pulse energy (>10¹³ photons/pulse) and linewidth (<10 GHz).
- Our setup:
 - Two cw Ti:Sa ring lasers as seed.
 - Pulsed dye amplifiers (~60 mJ/pulse, 30 Hz repetition rate).

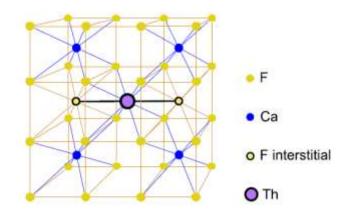




PTB - TU Wien cooperation: Laser excitation of ²²⁹Th-doped calciumfluoride crystals

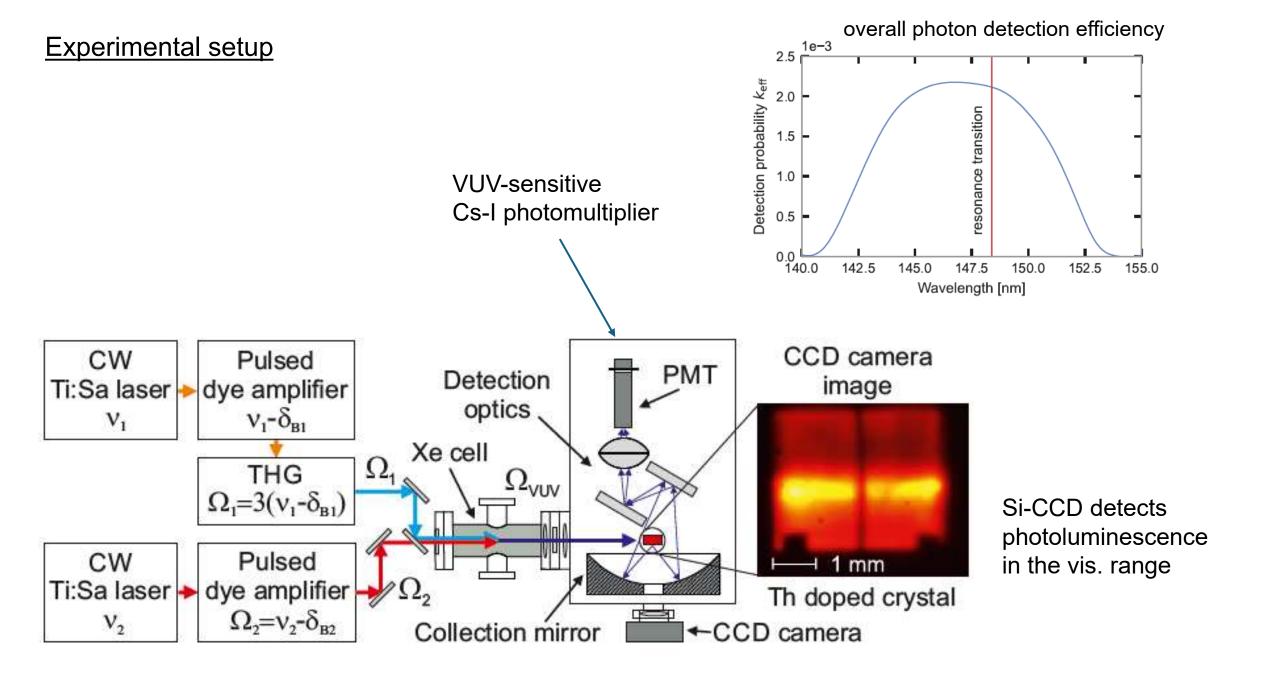




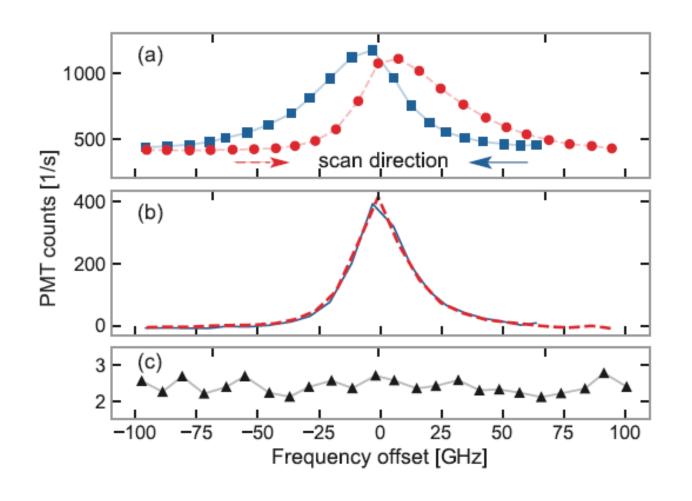








Resonant laser excitation, detected in VUV fluorescence



Excitation spectra at 148 nm: Hitting a narrow line with a broad laser.

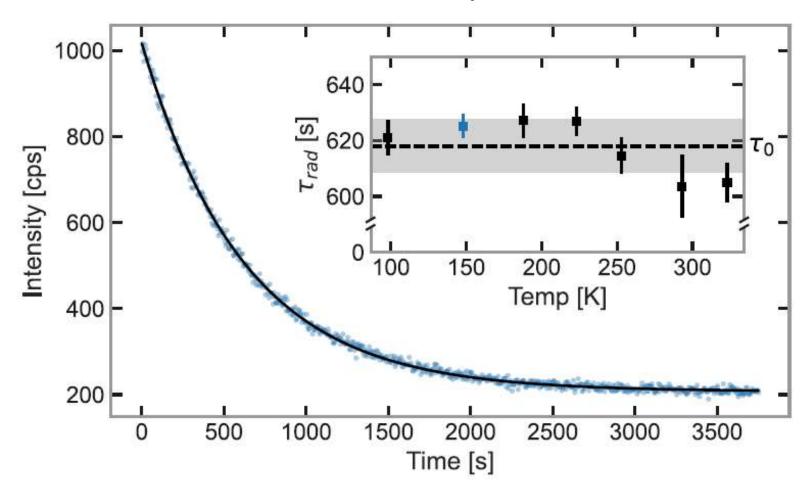
(each point: 120 s excitation, 150 s detection)

Line shapes after correction for the slow exponential fluorescence decay

Control experiment with Th-232: no signal

J. Tiedau et al. (PTB – TU Wien cooperation), Phys. Rev. Lett. **132**, 182501 (2024)

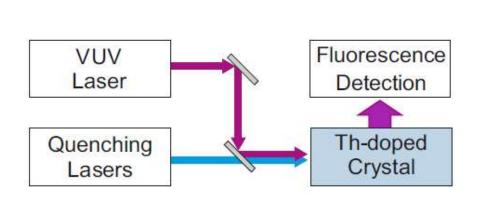
Fluorescence decay curves

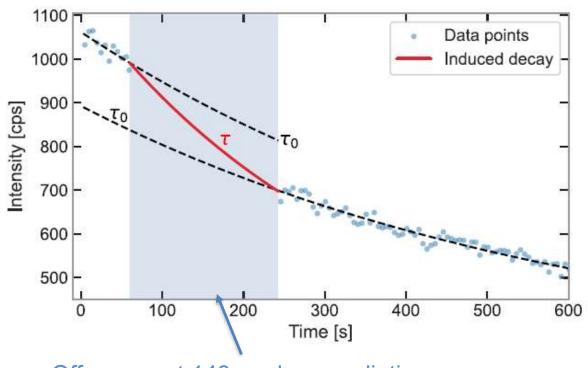


Decay time constant: 618(9) s:

- Identical for differently doped X2 and C10 crystals
- Identical for C10 before and after refluorination with CF₄
- Independent of crystal temperature 100 320 K
- Expected to be enhanced by the mode density in the crystal ∞ n³ \rightarrow isomer half-life 1740 s

Laser-Induced Quenching of the Th-229 Nuclear Clock Isomer in Calcium Fluoride

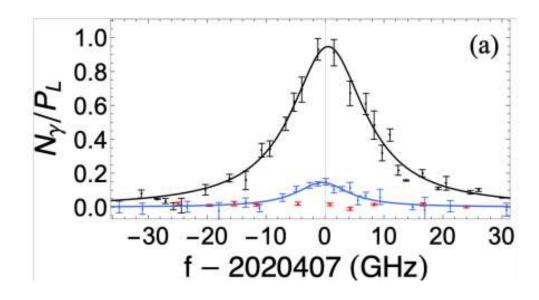




Off-resonant 148 nm laser radiation

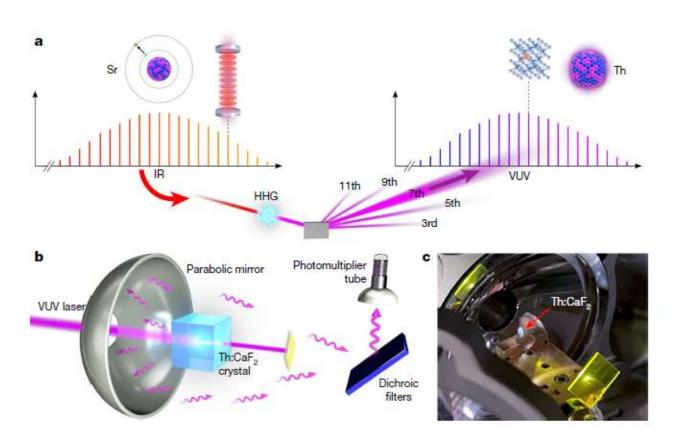
F. Schaden et al. (TU Wien – PTB), Phys. Rev. Res. **7**, L022036 (2025)

Th-229 excitation in Th:LiSAF at UCLA



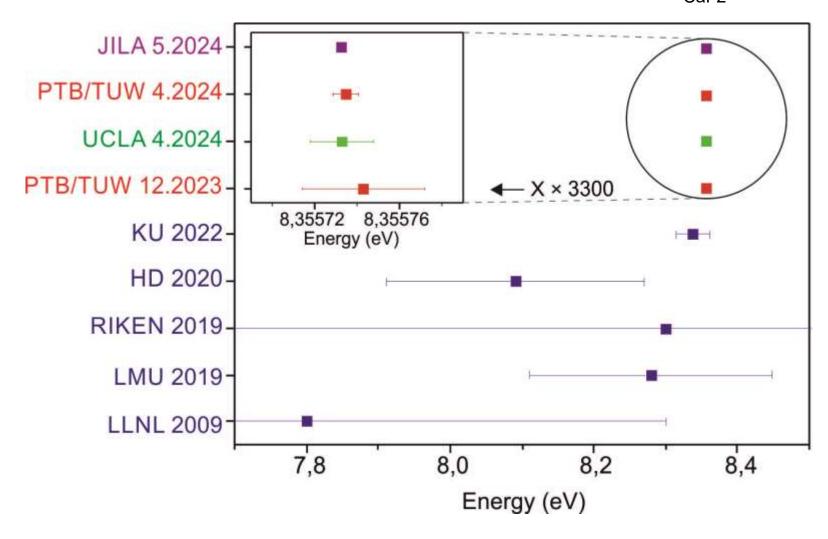
R. Elwell et al., Phys. Rev. Lett. 133, 013201 (2024)

Th-229 excitation with the 7th harmonic of a fs-laser



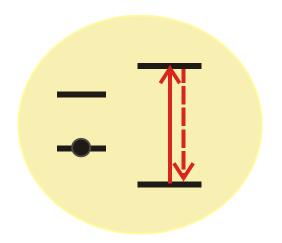
Ch. Zhang et al. (JILA and TU Wien), Nature **633**, 63 (2024)

 v_{CaF2} = 2020.407384335(2) THz

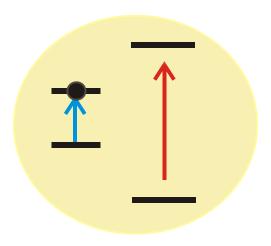


C. Zhang et al., Nature 633, 63 (2024)

<u>Trapped ions: Detection of the Nuclear Excitation in Nuclear-Electronic Double-Resonance</u>



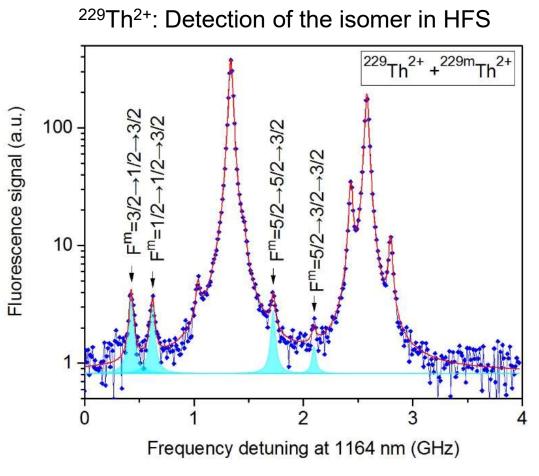
Nucleus in the ground state; laser-induced fluorescence from the shell.



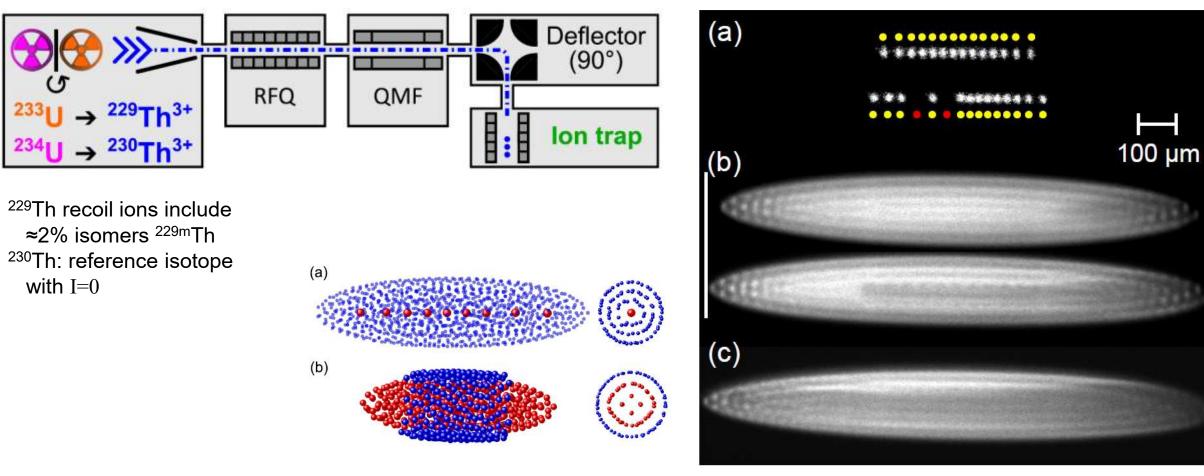
Laser excitation of the nucleus; change of hyperfine structure detected in intensity or polarisation of fluorescence.

Analog of Dehmelt's "electron shelving"
Observation of "quantum jumps" in the single-ion fluorescence
Coupled degrees of freedom and long nuclear coherence time

E. Peik, Chr. Tamm, Europhys. Lett. 61, 181 (2003)



J. Thielking et al. (PTB & LMU), Nature **556**, 321 (2018)

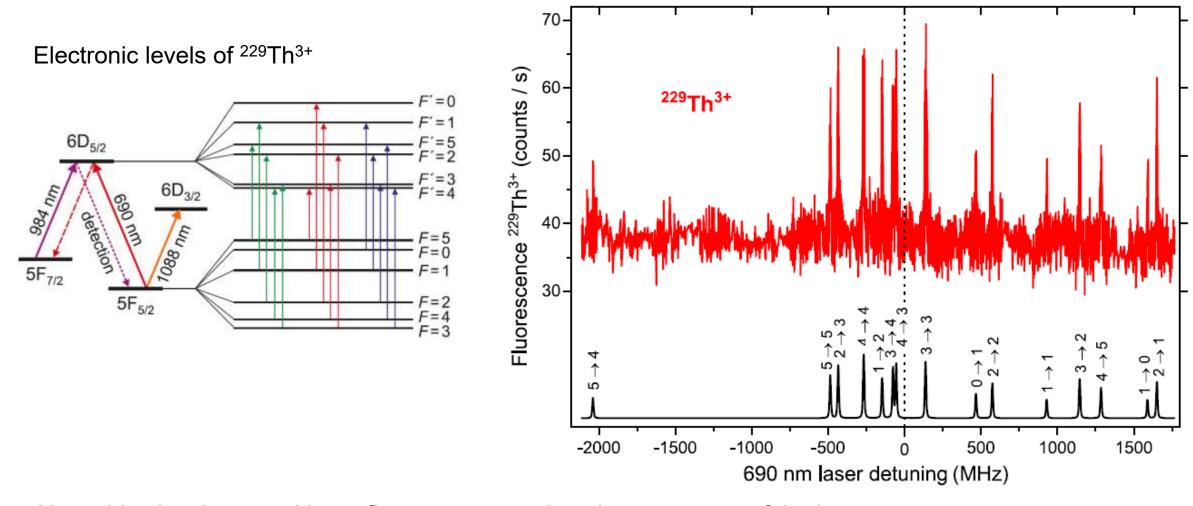


Simulated Coulomb crystals: blue Sr⁺, red: Th³⁺

Fluorescence of Sr⁺, Th³⁺ appear dark

G. Zitzer, J. Tiedau, M. V. Okhapkin, K. Zhang, C. Mokry, J. Runke, Ch. E. Düllmann, E. Peik, Phys. Rev. A **109**, 033116 (2024)

Hyperfine structure of ²²⁹Th³⁺ in the nuclear ground state



Next objective: improved hyperfine structure and nuclear moments of the isomer.

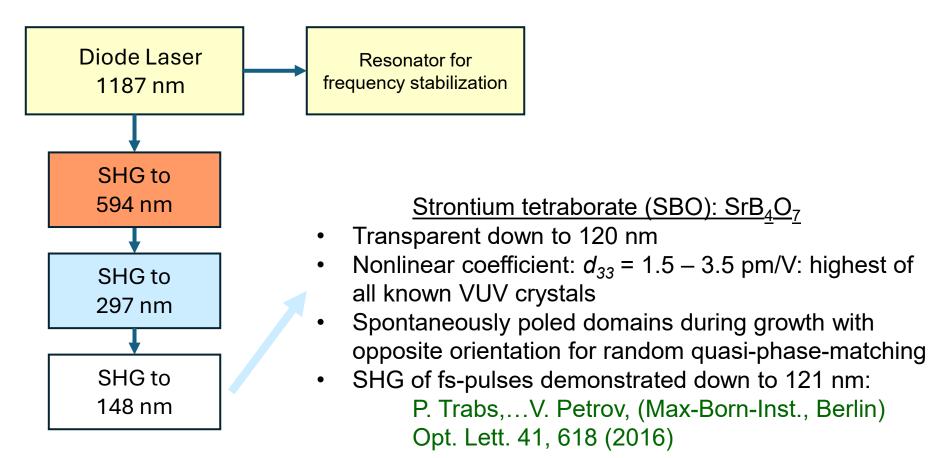
→ relevant info on nuclear charge distribution and nuclear transition matrix element.

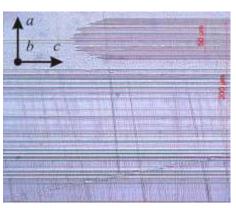
G. Zitzer, J. Tiedau, Ch. E. Düllmann, M. V. Okhapkin, E. Peik, Phys. Rev. A 111, L050802 (2025)

A narrow-linewidth CW laser source for 148 nm



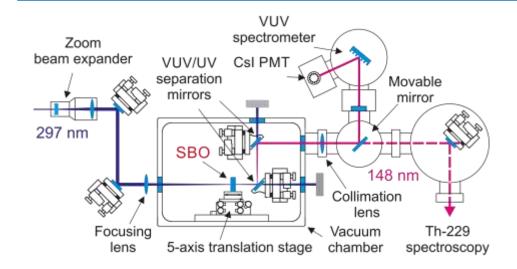
- Infrared diode laser at 1187 nm, stabilized to reference resonator (e.g. cryo silicon)
 → <<1 Hz linewidth demonstrated
- VUV generation by 3× successive frequency doubling in nonlinear optical crystals

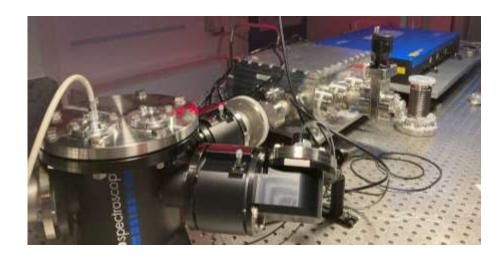




148 nm CW laser source



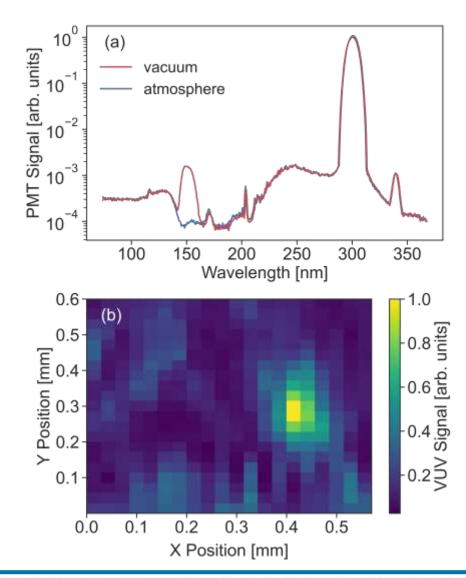




- CW radiation at 297 nm from frequency-quadrupled laser (Toptica TA-FHG pro)
- SBO is mounted on a 5-axis translation stage.
- The main challenge of SHG detection is to achieve sufficient suppression of fundamental radiation.
- To detect VUV signal in fW range from ~0.3 W fundamental radiation, a suppression factor of 10¹⁴ is required.
- Two dichroic mirrors in combination with a spectrometer grating, and solar-blind CsI PMT with low UV sensitivity.

SHG optimization

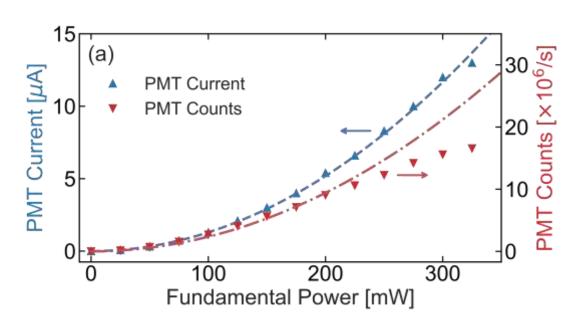




- The recorded PMT signal showing fundamental and SHG spectra.
- The SHG signal at 148 nm vanishes under atmospheric condition due to high absorption of the VUV light

- The irregular domain structure of SBO crystal is inhomogeneous across the sample.
- A beam radius of ~ 25 μm for the fundamental radiation was used.





- Signal detection: photon counting and PMT current measurements.
- Both PMT current and PMT count signals show quadratic dependence on the fundamental power below 200 mW.
- The generated VUV power is 1.3^{+0.7}_{-0.5} nW at an incident UV power of 325 mW.
- The large uncertainty arises from discrepancies between the two-detection method and the uncertainty in the detection efficiency of the setup.

V. Lal, M. V. Okhapkin, J. Tiedau, N. Irwin, V. Petrov, E. Peik, arXiv:2507.17719

See also: Qi Xiao et al., arXiv:2507.19449: CW four-wave-mixing in cadmium vapor

What's next?

- Building narrow-linewidth VUV lasers
- Dynamics and lineshapes in the Th-doped crystals
- New materials, different host crystals
- Exciting the nuclear transitions in trapped Th ions
- Th-229 nuclear clocks in tests of fundamental physics

PTB Team:

J. Tiedau

M. Okhapkin

K. Zhang

J. Thielking

G. Zitzer

V. Lal

N. Irwin

S. Maurya

E. Peik

MBI:

V. Petrov

TU Wien Team:



T. Pronebner

I. Morawetz

L. Toscani De Col

F. Schneider

A. Leitner

M. Pressler

G. Kazakov

K. Beeks

T. Sikorsky

T. Schumm









